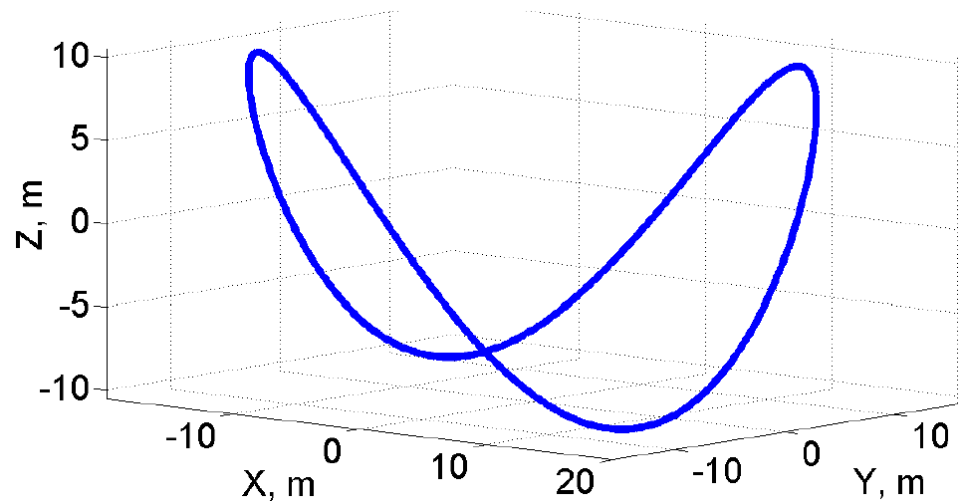
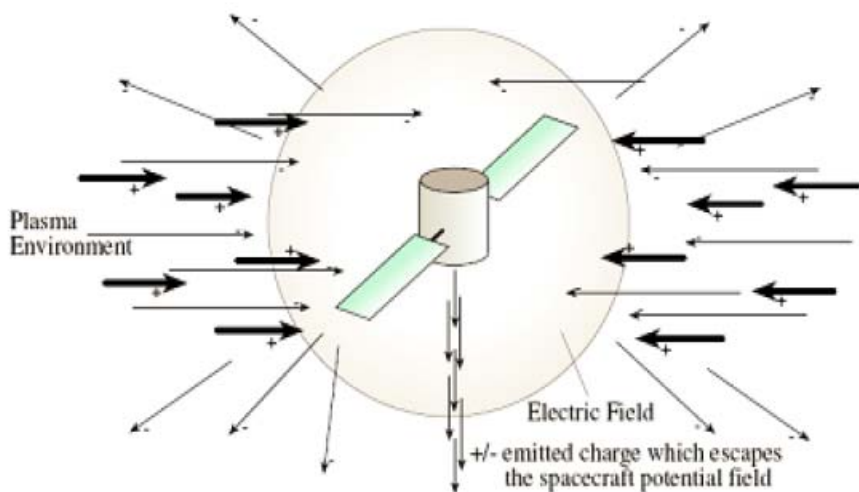


Periodic Relative Orbits of Two Spacecraft Subject to Differential Gravity and Coulomb Forces

Presented by Daan Stevenson for
Drew R. Jones (author) & Hanspeter Schaub (co-author)

5th SFFMT Conference

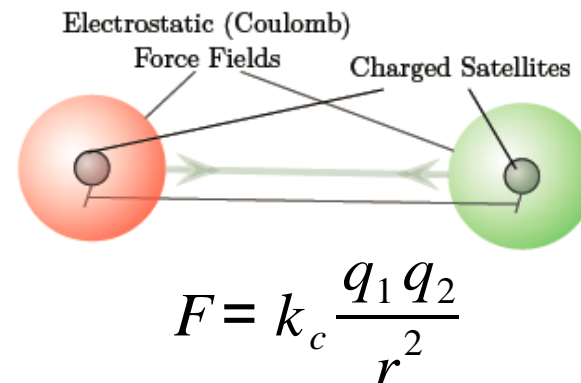
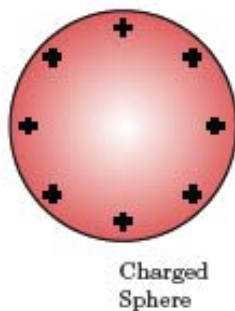
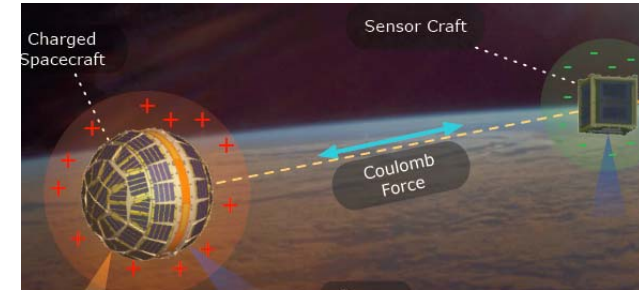
May 31, 2013



Coulomb Formation Motivation

Coulomb Formation: Close-flying charged craft using electrostatic forces

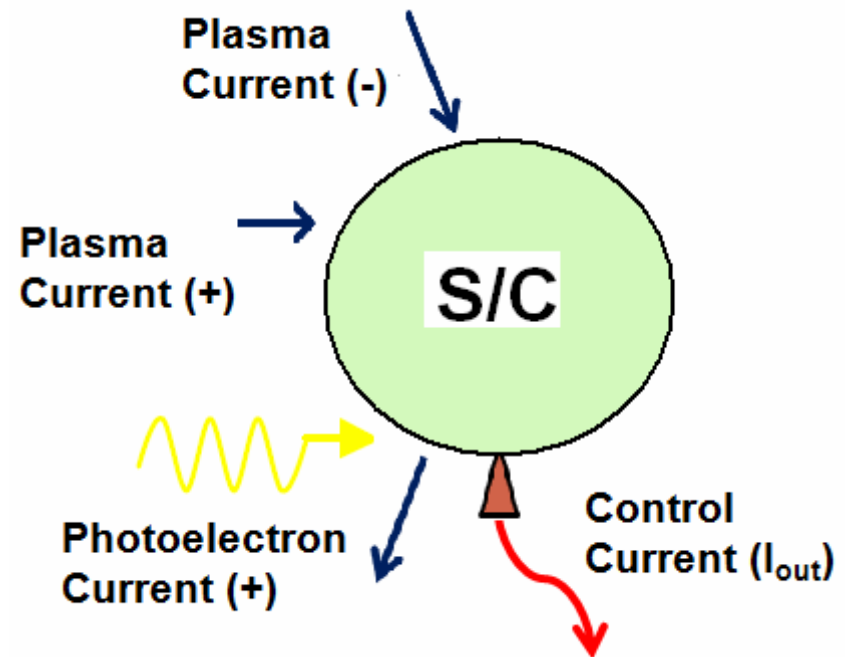
- Very efficient (ISP~ 10^{13} s)
- Forces have limited reachability
- Forces introduce NL coupling
- Forces complex to model in the plasma



- I. Approximate s/c charging and Coulomb force model
- II. Dynamics and dynamical system properties
- III. Derivation of periodic Coulomb formation solutions
- IV. Simulated periodic Coulomb formation solutions
- V. Perturbed Coulomb formation motions
- VI. Conclusions and future research

Spacecraft Charging

- Spacecraft naturally assume non-zero potential (charge)
- Artificially altering potential has been demonstrated
- Spherical s/c model adopted
 - Somewhat an abstraction
 - Enables 1st order calculations and assumptions

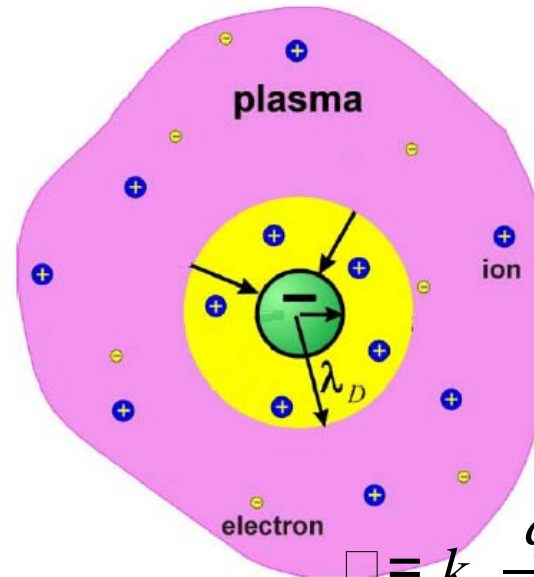


Electrostatic Model

- Spacecraft are not point charges in a vacuum...
 - Adopt approximate and analytical Coulomb force model
 - But still account for plasma shielding effect
 - Assume formations near GEO

Coulomb Force: Point Charges in Vacuum

$$F_{12} = k_c \frac{q_1 q_2}{r_{12}^2}$$



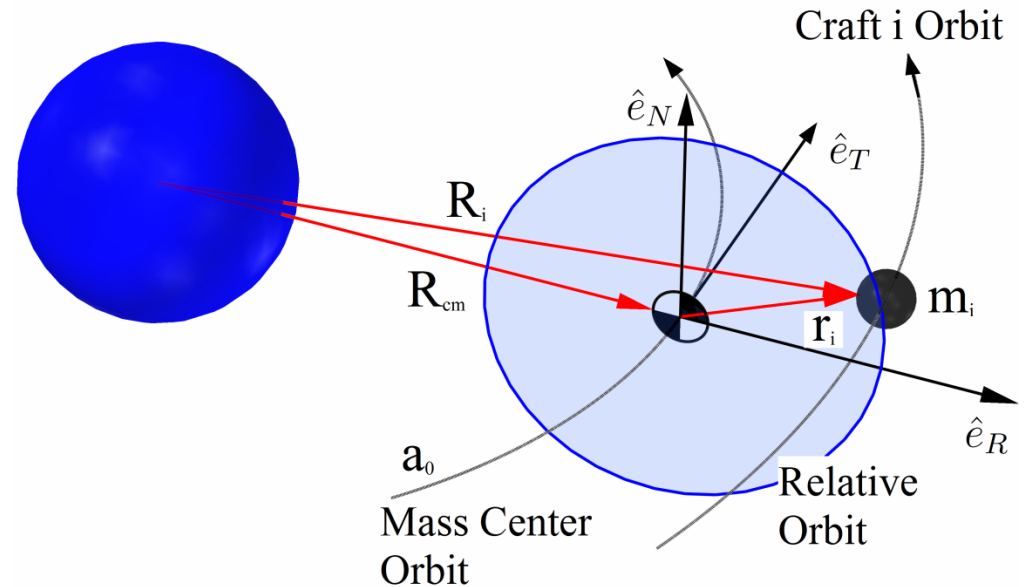
Finite Bodies with Plasma Shielding

$$F_{12} = \frac{q_1 q_2 (1 + r_{12}/\lambda_d)}{r_{12}^2 \exp[r_{12}/\lambda_d]}$$

$$E_i = k_c \frac{q_i}{R_{sc}}$$

- Hill Frame Model
- Clohessy-Wiltshire relative motion EOM

$$\sum_i m_i \mathbf{r}_i = 0$$



- With Debye-Hückel Coulomb force model:

$$\ddot{\mathbf{r}}_i = \begin{bmatrix} 2\omega\dot{y}_i + 3\omega^2 x_i \\ -2\omega\dot{x}_i \\ -\omega^2 z_i \end{bmatrix} + \left[\frac{k_c q_i}{m_i} \sum_{j \neq i} \frac{q_j e^{-r_{ij}/\lambda_d}}{r_{ij}^3} \left(1 + \frac{r_{ij}}{\lambda_d} \right) \mathbf{r}_{ij} \right]$$

- Dynamics normalized, one craft removed
 - s/c 2 motion is always dependent, via CM constraint
 - Scaled charge products introduce time transformation

$$Q_{12} = q_1 q_2$$
$$\tilde{Q}_{12} = \frac{k_c Q_{12}}{\omega^2} \quad d\tau = \omega dt \quad (\zeta)' = \frac{d\zeta}{d\tau} = \frac{1}{\omega} \frac{d\zeta}{dt}$$

- Scalar constant of motion is shown to exist
 - Holds for N-craft Hill frame systems with internal forcing
 - For two-craft formation, the motion constant is:

$$x(\tau)y''(\tau) - y(\tau)x''(\tau) = 0$$

- Seek periodic solutions to vector ODE $\mathbf{X}' = \mathbf{F}(\mathbf{X}, \mathbf{u}, \tau)$
 - Dynamics linearized about such solutions

$$\delta \mathbf{X}'(\tau) = \left. \left(\frac{\partial \mathbf{F}}{\partial \mathbf{X}} \right) \right|_{(\mathbf{X}^*, \mathbf{u}^*)} \delta \mathbf{X}(\tau) = \mathbf{A}(\tau) \delta \mathbf{X}(\tau) \quad \delta \mathbf{X}(\tau) = \Phi(\tau, 0) \delta \mathbf{X}(0)$$

$$\Phi'(\tau, 0) = \mathbf{A}(\tau) \Phi(\tau, 0)$$

- STM matrix has identical symplectic property as known for linear dynamics about libration points in CRTBP
- Asymptotically stable solutions cannot exist

$$\mathbf{J}\mathbf{A}^T = -\mathbf{A}\mathbf{J} \quad \Phi\mathbf{J}\Phi^T = \mathbf{J} \quad \mathbf{J} = \left[\begin{array}{c|c} \mathbf{0} & \mathbf{I} \\ \hline -\mathbf{I} & \mathbf{G} \end{array} \right]$$

1. $\det(\Phi) = |\Phi| = 1$

2. At least one Floquet multiplier has modulus of unity: $|\sigma_i| = 1$

3. The σ_i appear in reciprocal pairs (i.e. if σ_i is eigenvalue, then so is $\sigma_j = 1/\sigma_i$)

Periodic Coulomb Formations

- Assume positions as simple harmonic oscillators
 - Solve time-varying charges that produce assumed motion
 - Motions are restricted (cannot be assumed arbitrarily)
- Relative instability measured via Floquet multipliers
- Dynamical coupling permits 3 solution families
 - Orbit-Normal oscillations only (1D)
 - Reference orbit-plane (2D) motions
 - Full state (3D) motions

Reference Orbit Plane Solutions

$$x(\tau) = A_x \cos(\theta_x \tau) \quad y(\tau) = A_y \sin(\theta_y \tau)$$

$$1) \quad x(\tau)y''(\tau) - y(\tau)x''(\tau) = 0$$



$$\theta_x = \theta_y$$

$$2) \quad \text{Dynamical Coupling} \\ \text{(Coulomb terms):}$$

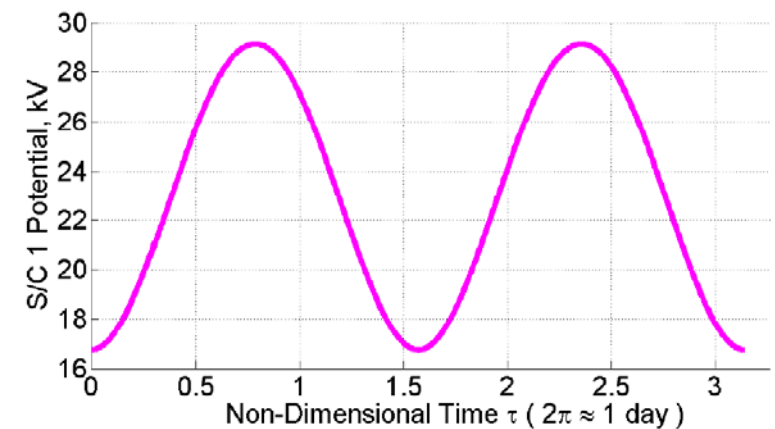


$$\left(\frac{A_y}{A_x} \right) = \frac{-3 \pm \sqrt{9 + 16\theta^2}}{4\theta}$$

- Two roots of quadratic yield two families (cases A/B)

- Relative orbit is ellipse
- Major axis along either radial (case A) or transverse (case B)
- Short period (case B) orbits → least unstable

$$\tilde{Q}(r(\tau)) = \frac{-1}{\Psi(r)} \left[\theta^2 + 3 + \left(\frac{-3 \pm \sqrt{9 + 16\theta^2}}{2} \right) \right]$$



(b) Case B: $A_x < A_y$

Full State Solutions

- Same in-plane assumed functions (same restrictions)

- Additionally

$$z(\tau) = A_z \sin(B_z \theta \tau)$$

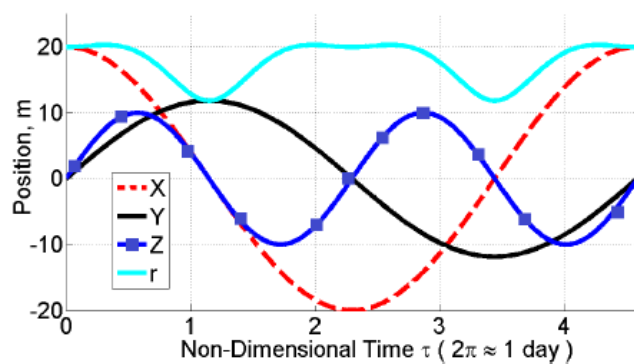
- 1) **Out of plane frequency factor must be even integer**

$$\longrightarrow B_z = 2, 4, 8 \dots$$

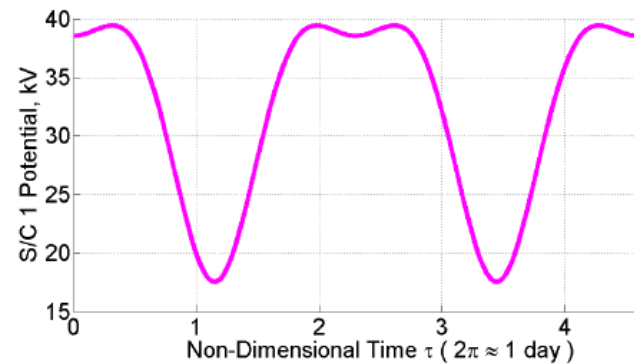
- 2) **Time period no longer free, must satisfy:**

$$\longrightarrow 8\theta^2 + \left(-3 \pm \sqrt{9 + 16\theta^2}\right) [\theta^2(1 - B_z^2) + 1] = 0$$

- Projected trajectory is ellipse aligned as in case A/B
- Case B orbits have smaller max. Floquet multiplier

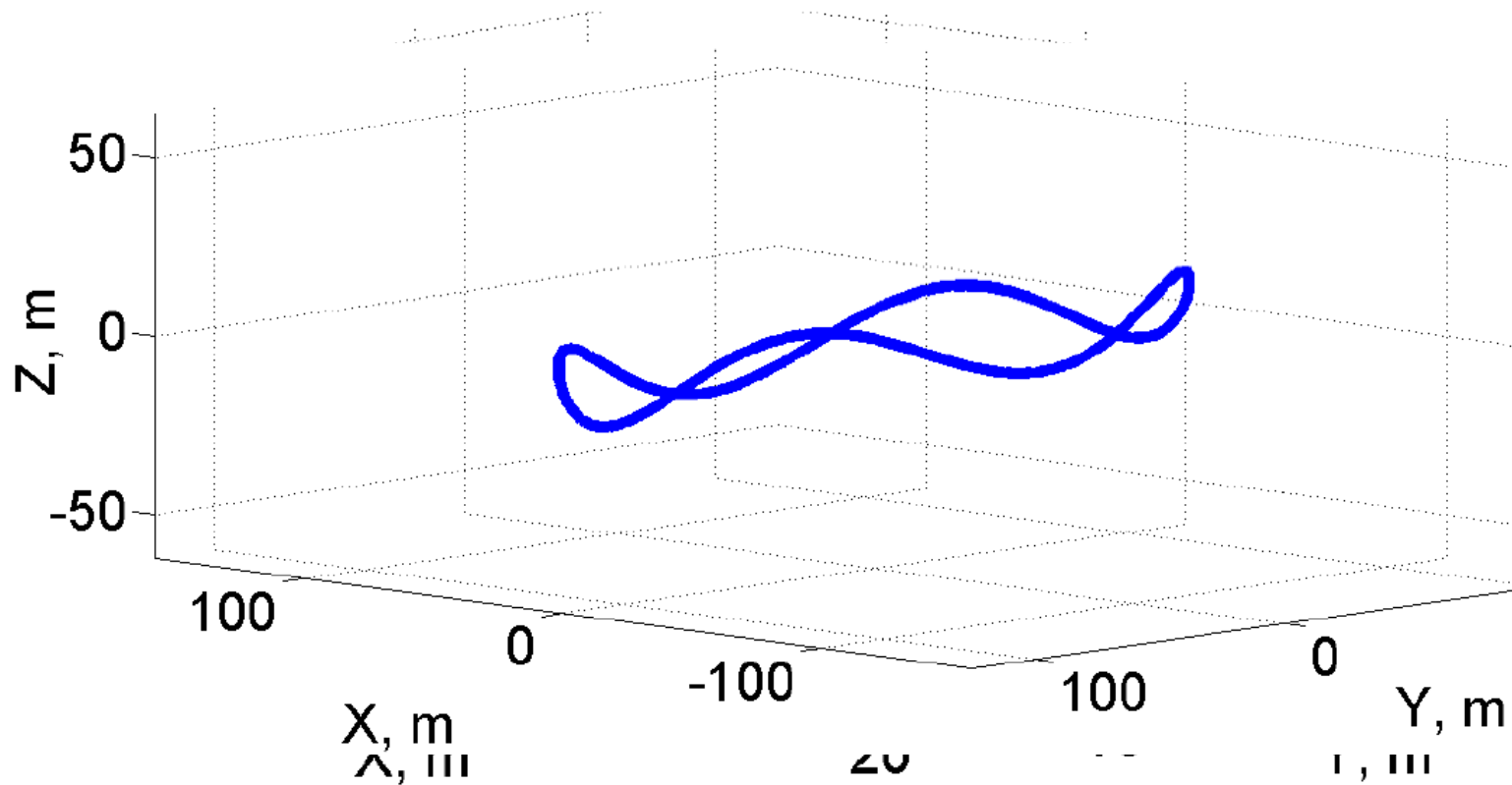


(a) S/C 1 Position History



(b) S/C 1 Potential History

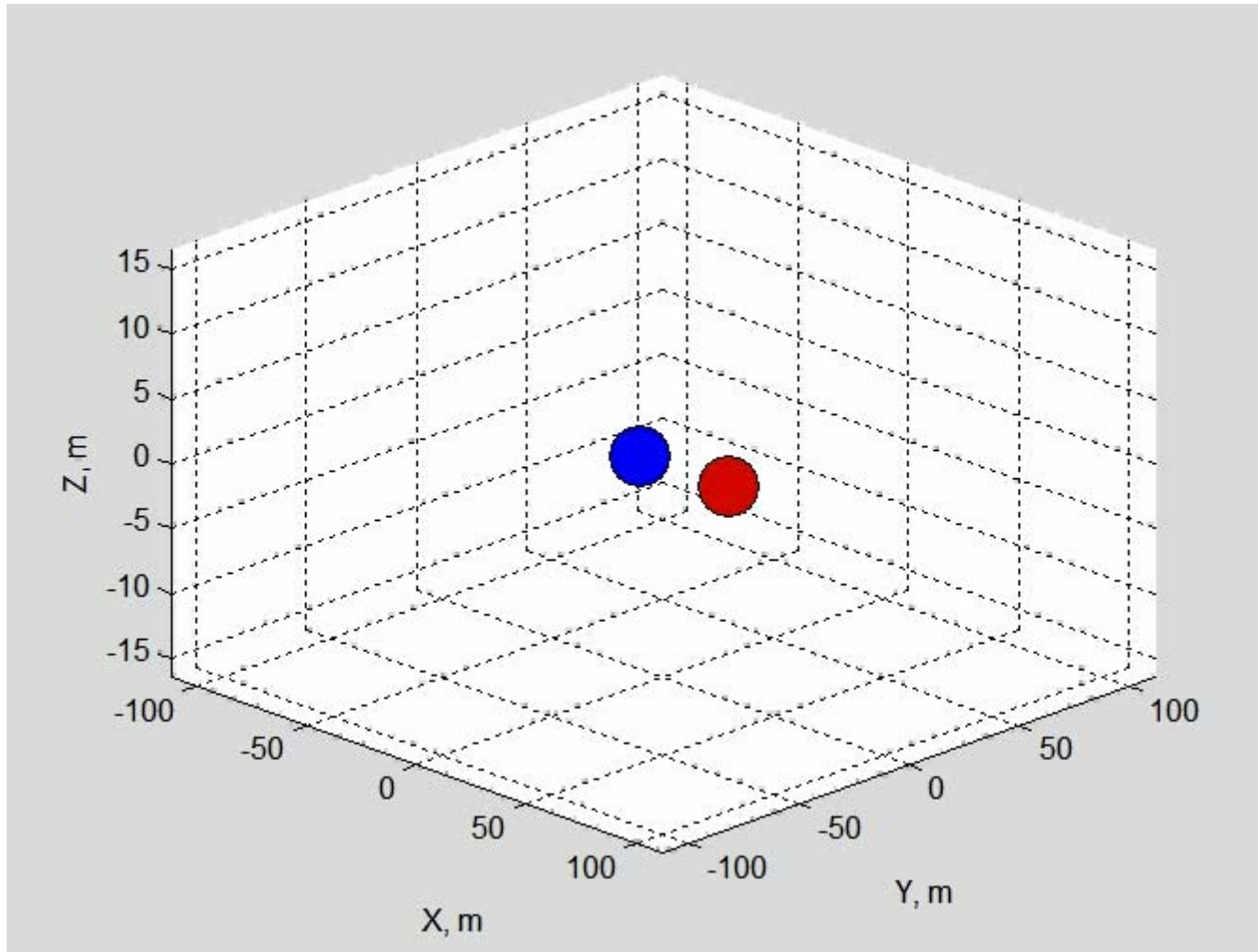
- Example relative orbit
 - S/C 1 Trajectory \rightarrow like saddle, x-y projection is ellipse



Full State Periodic Solution: Case B, $B_z=2$

Circumnavigating 2-s/c Orbit

- Another Example
 - Both craft trajectories shown, different mass vehicles



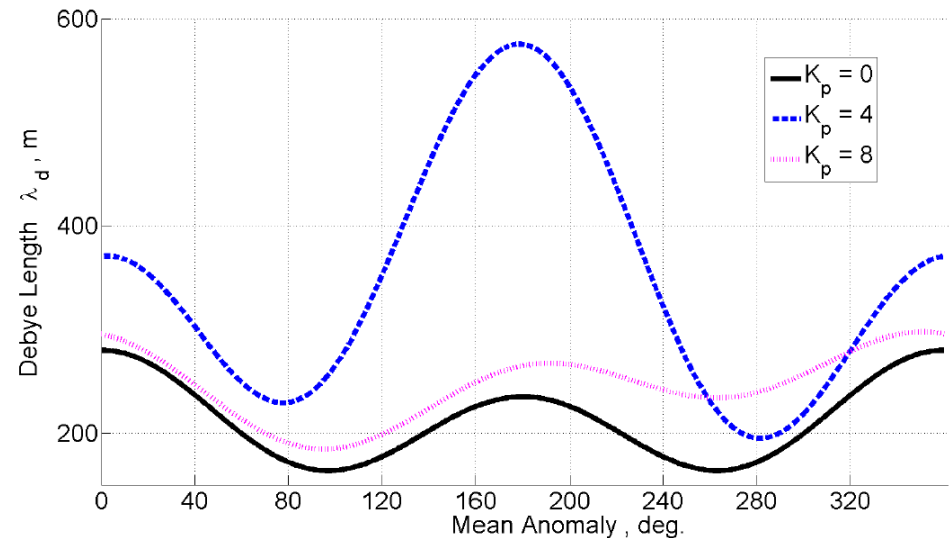
Perturbed Periodic Coulomb Formation Solutions Propagated in Inertial Frame

- Accuracy of solutions in higher fidelity model
 - Revert back to Newtonian gravity, inertial frame
 - Transformations between ECI and Hill frames
 - Include primary perturbations

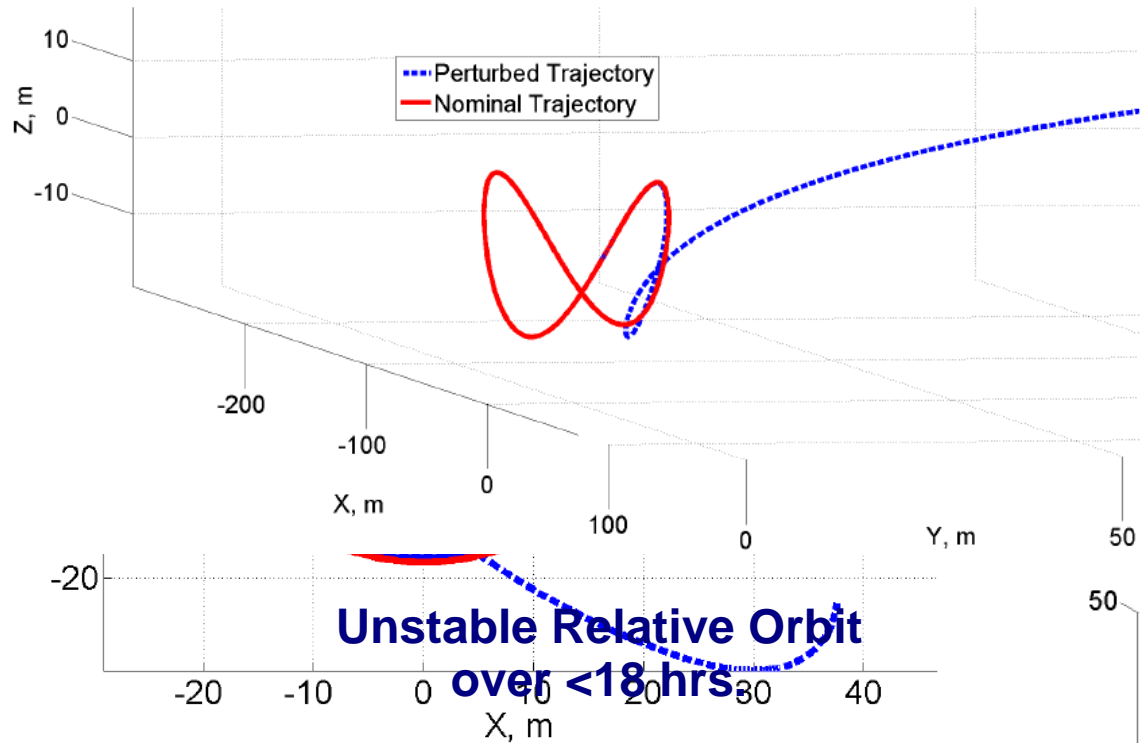
Solar Radiation Pressure using “Cannonball Model”

$$f_{\text{srp}} = C_R \frac{\pi R_{sc}^2 \Theta}{m_i c}$$

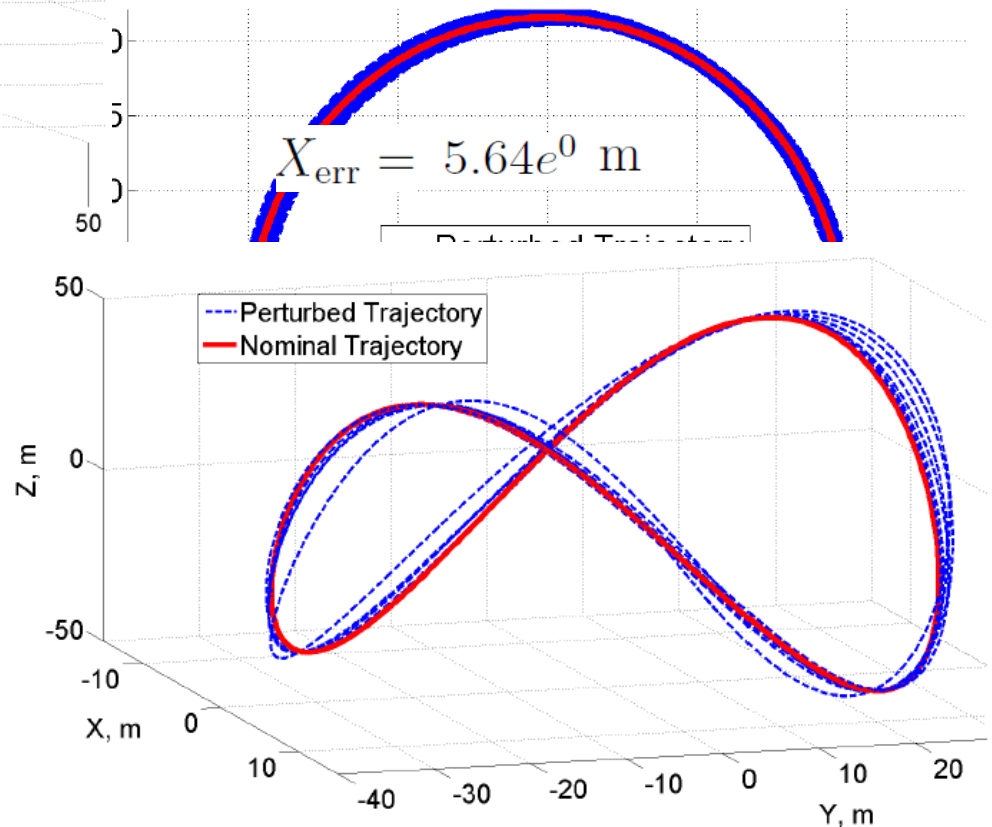
Induced Perturbation due to Parametric Uncertainty (Model Error) in Debye Length



Perturbed Periodic Solutions



Unstable Relative Orbit over 10 Revolutions



- Application of new theory opens multiple pathways
 - The existence of periodic solutions enabled by open-loop electrostatic forcing
 - Stability properties assessed via Floquet theory
 - Some periodic solutions are very nearly stable
- Future Work
 - Other admissible periodic solutions (represent w/ finite Fourier)
 - Feedback control of periodic solutions
 - Test these solutions (and derive new ones) with relaxed assumptions and higher order Coulomb force modeling
 - Precisely measure, control, and estimate s/c potentials (and Coulomb forces) in order to realize these motions in practice?

QUESTIONS

Please refer further questions to Drew Jones: drjones604@gmail.com

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EXTRA SLIDES

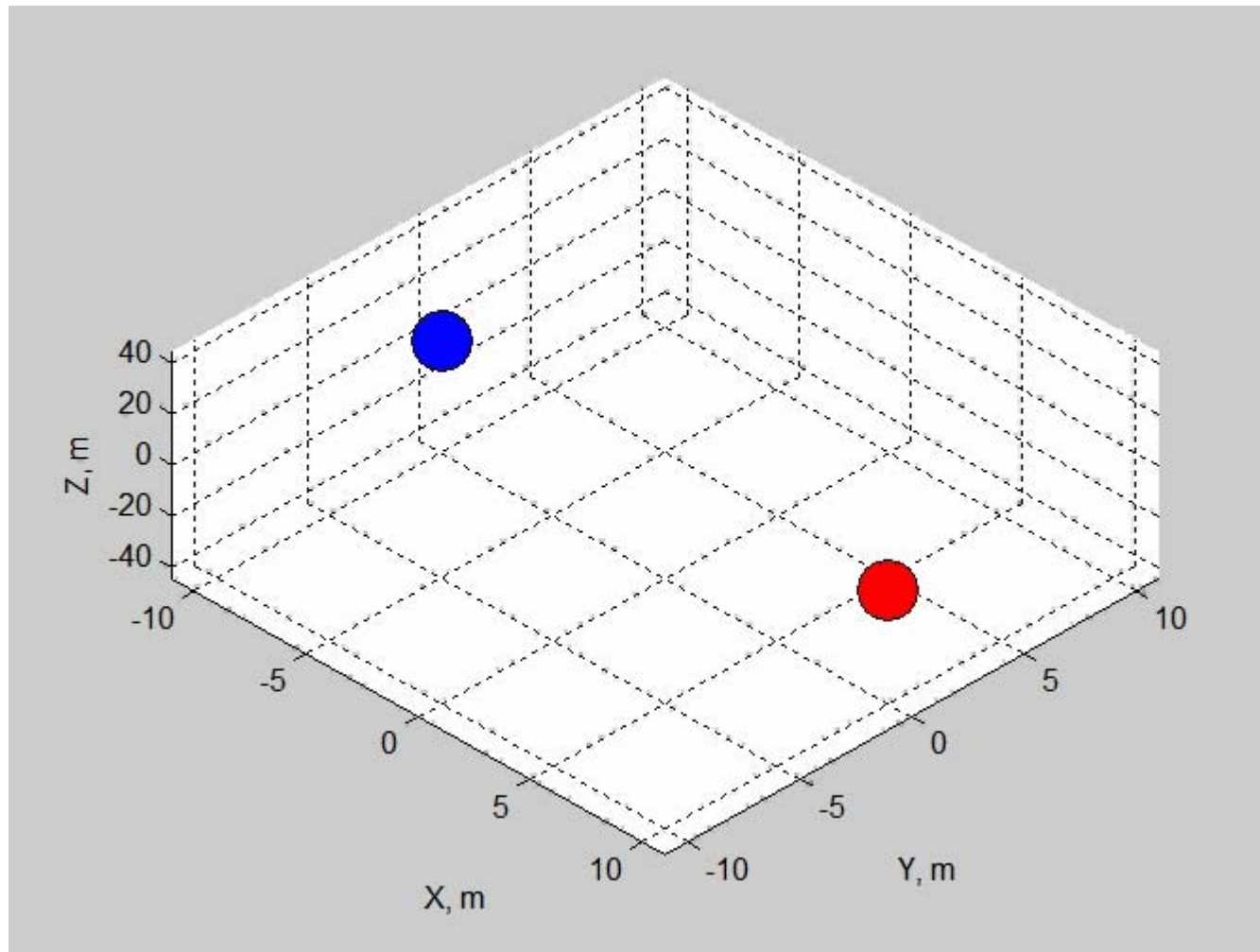
- True force between finite s/c: solution to time-dependent Vlasov-Poisson PDE

$$F \approx k_c \frac{q_1 q_2 (1 + r_{12}/\lambda_d)}{r_{12}^2 \exp[r_{12}/\lambda_d]} \quad \square_i = k_c \frac{q_i}{R_{sc}}$$

- Assumptions for approximate model accuracy:
 - Craft act equivalently as spheres at a distance
 - Conservative plasma shielding
 - Decoupled capacitances
 - Constant and finite λ_d (at GEO order of 100 m)
 - Formations near GEO (λ_d negligible over R_{sc})
 - Lower bound on separation distance: $r_{12} > 10 R_{sc}$

Periodic Coulomb Formation

- Case A, $B_z = 2$, $A_x = 10$ m, $A_y = 6$ m, $A_z = 40$ m (17 hrs)



Periodic Coulomb Formation

- Case B, $B_z = 4$, $A_x = 20$ m, $A_y = 125$ m, $A_z = 5$ m (4 days)

